

## Evaluation of moments and their application in Hückel molecular orbital theory

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Moments of acyclic carbon chains are treated in a systematic way. Explicit formulas in terms of connectivities are tabulated up to  $\mu_{14}$ . This facilitates the evaluation of moments by simply counting the numbers of various fragments involved. The total  $\pi$  electron energy is analyzed by means of moments and the meaning of additivity is interpreted. An approximate formula for  $E_\pi$  is parametrized by truncation, preserving only five bond parameters. Based on these, we attempt to better rationalize and reformulate the concept of aromaticity.

**Key words:** Hückel theory—Graph theory—moments—aromaticity

### 1. Introduction

The traditional way for dealing with MO (abbreviation of “molecular orbital”) calculations is to solve a secular equation [1]

$$|ES - H| = 0. \quad (1)$$

This equation depends on the basis set  $\{\phi_i\}$  selected, because the overlap and Hamiltonian matrices  $S$  and  $H$  are defined in terms of  $\{\phi_i\}$

$$S = (S_{ij}) \quad S_{ij} = \langle \phi_i | \phi_j \rangle \quad (2)$$

$$H = (H_{ij}) \quad H_{ij} = \langle \phi_i | \hat{H} | \phi_j \rangle.$$

\* Formerly used names: Kiang, Yuan-sun; Tang, Au-chin

Usually, energy levels are obtained first, and wave functions and other observable calculations follow. If  $S$  is approximated by the unit matrix

$$S = I, \quad S_{ij} = \delta_{ij} \quad (3)$$

then Eq. (1) is easily expanded into a polynomial characterized by powers of  $E$  [2-4]

$$E^N + a_1 E^{N-1} + a_2 E^{N-2} + \cdots + a_N = 0. \quad (4)$$

The coefficients  $a_1, a_2, \dots, a_N$  are related to the energy eigenvalues  $E_1, E_2, \dots, E_N$ , thus they can be equally well utilized for handling MO problems. Still other possibilities exist. Among them, moments are a useful analytical tool. They are defined as

$$\mu_l = E_1^l + E_2^l + \cdots + E_N^l \quad (l=0, 1, 2, \dots) \quad (5)$$

and are equal to the trace of  $l$ -th power of the matrix  $H$

$$\mu_l = \text{trace } H^l \quad (6)$$

Obviously, only the lower  $N$  members constitute an independent set, since all of the higher moments with  $l \geq N$  are equal to a linear combination of preceding  $N$  members, according to Eq. (4). Namely

$$\mu_{N+m} = a_1 \mu_{N+m-1} + a_2 \mu_{N+m-2} + \cdots + a_N \mu_m \quad (m=0, 1, 2, \dots). \quad (7)$$

The importance and advantage of utilizing moments rests upon their topological meaning in relationship to the connectivity of a molecule. The situation is particularly simple and elegant for the Hückel theory. After putting

$$\alpha = 0, \quad \beta = 1. \quad (8)$$

$H$  becomes an adjacency matrix  $A$ , with its entries satisfying

$$A_{ij} = 1 \quad \text{when } i, j \text{ connected} \\ = 0 \quad \text{otherwise.}$$

In this case, each non-zero term in  $\mu_l$  is exactly equal to unity

$$\mu_l = \sum_{i=1}^N (A^l)_{ii} = \sum_{i=1}^N \sum_{\alpha, \beta, \dots, \delta} A_{i\alpha} A_{\alpha\beta} \cdots A_{\delta i} = \sum_{i=1}^N \sum_{\alpha, \beta, \dots, \delta} (1) \quad (10)$$

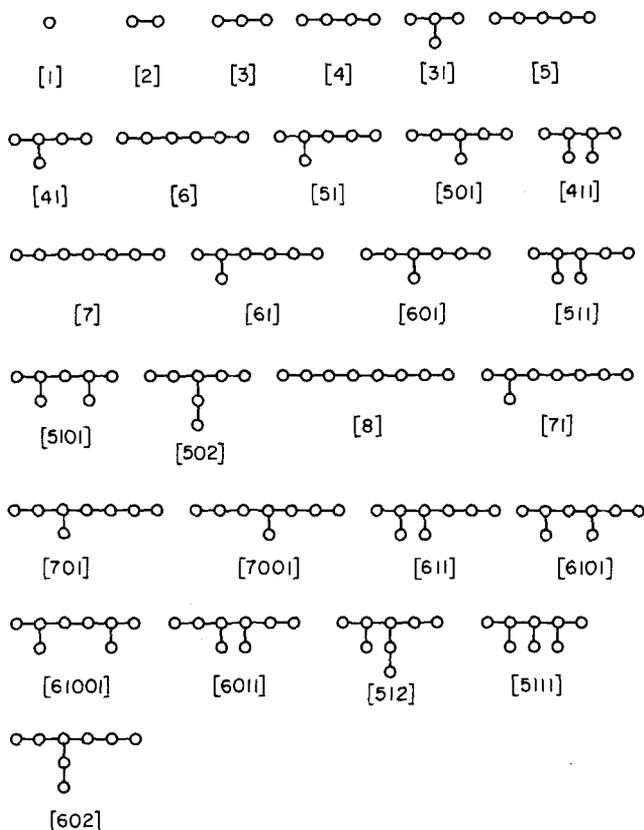
representing a cyclic walk of length  $l$  starting from vertex  $i$  and passing through vertices  $\alpha, \beta, \dots, \delta$  irrespective of the repeated appearance of some vertices. The sum of such walks with respect to each vertex is the value of  $\mu_l$ , namely the number of random closed walks of length  $l$ .

Directly enumerating moments is feasible for small molecules but becomes tedious when the molecule is large. In this paper, we present a procedure for evaluating moments and list formulas for the convenience of enumerating the moments of acyclic chains. It is proposed that the total  $\pi$  electron energy be expanded in terms of moments. A better interpretation of additivity may thus be achieved.

We reformulate a previous theory of aromaticity by parametrization of the total energy in terms of five branching types of bonds.

## 2. Moments of acyclic systems

We limit our attention to acyclic conjugated skeletons with a maximum atom valency of three. Nevertheless, the method proposed is general. First of all, it seems worthwhile to discuss how to characterize these skeletons or fragments. Let us present the smaller members up to  $N = 8$  by diagrams and symbols



The symbol  $[N - n; a, b, c, \dots, n - a - b - c - \dots]$  temporarily used here signifies the topology of the fragment,  $N - n$  representing the length of the main chain and  $a, b, c, \dots$  signifying lengths of the first, second, ... straight side chains respectively. For higher members, where branching side chains exist, this symbol should be modified, but we omit doing that for simplicity. These diagrams can represent real molecules as their carbon skeletons or they can represent fragments that a real molecule includes.

Because an acyclic chain is alternant, with vanishing odd moments, we concentrate on even moments only. Obviously for any molecule  $G$  we have its moments

satisfying

$$\mu_{2n}^G = \sum_{G'} C_{2n}^{G'} N_{G'} \quad (11)$$

where  $G'$  represents a fragment of  $G$  involving bonds equal to or less than  $n$ ,  $N_{G'}$  signifies the appearances of  $G'$  and  $C_{2n}^{G'}$  counts the random closed walks of length  $2n$  spanned by  $G'$ . All of these are well illustrated by a simple example shown below.

$G =$  butadiene

$n = 3$	possible $G'$	[4]	[3]	[2]
	$N_{G'}$	1	2	3
$n = 2$	$G'$	[3]	[2]	
	$N_{G'}$	2	3	

An expression for the coefficient  $C_{2n}^G$  in terms of  $\mu_{2n}^{G'}$  can be derived from Eq. (10) by induction. It takes the following form

$$C_{2n}^G = \mu_{2n}^G - \sum_t N_{G-t} \mu_{2n}^{G-t} + \sum_{t,t'} N_{G-t-t'} \mu_{2n}^{G-t-t'} - \sum_{t,t',t''} N_{G-t-t'-t''} \mu_{2n}^{G-t-t'-t''} + \dots \quad (12)$$

where  $t, t', t'', \dots$  represent terminal atoms in molecule  $G$ ;  $G-t, G-t-t', \dots$  mean fragments generated by deleting one, two, ... or more terminals from  $G$  respectively. If we still take butadiene as an example, Eq. (12) is reduced to

$$C_{2n}^{[4]} = \mu_{2n}^{[4]} - 2\mu_{2n}^{[3]} + \mu_{2n}^{[2]} \quad (13)$$

### 3. Analytical $C_{2n}^G$ 's

Generally speaking, coefficients  $C_{2n}^G$  are too complicated to be given analytically except for some special  $G$  or special  $n$ . If no more than three terms are involved in the characteristic polynomial of  $G$ , namely

$$P_G = X^m(X^4 - a_1X^2 + a_2) \quad (14)$$

then by utilizing Eq. (7) recursively, an analytical form for  $\mu_{2n}^G$  will be obtained:

$$\mu_{2n}^G = g_{n-1}(a_1, a_2)\mu_2^G - g_{n-2}(a_1, a_2)a_2\mu_0^G \quad (15)$$

where

$$g_{n-1}(a_1, a_2) = \sum_r \binom{n-r}{r} (-1)^r a_1^{n-2r} a_2^r \quad (16)$$

The symbol  $\binom{n-r}{r}$  means the binomial coefficient and  $g_n(a_1, a_2)$ , a polynomial of  $a_1$  and  $a_2$ , is further reduced to a Tchebiecheff polynomial when  $a_2$  equals unity. In Eq. (15),  $\mu_2^G$  and  $\mu_0^G$  are simply analytical (see Eq. (19)). As a corollary, analytical  $C_{2n}^G$  are also deduced from Eq. (12). Some of them are listed below

$$\begin{aligned}
C_{2n}^{[0]} &= 1 \\
C_{2n}^{[2]} &= 2 \\
C_{2n}^{[3]} &= 2^{n+1} - 2^2 \\
C_{2n}^{[4]} &= 6g_{n-1}(3, 1) - 4g_{n-2}(3, 1) - 2(2^n - 1) \\
C_{2n}^{[31]} &= 6(3^{n-1} - 2^n + 1) \\
&\dots
\end{aligned}
\tag{17}$$

If the fragment  $G$  has its number of bonds equal to  $n$ , the coefficient  $C_{2n}^G$  is also given analytically, namely

$$C_{2n}^G = (2n)2^{N_3} \tag{18}$$

where  $N_3$  is the number of trivalent atoms. Eq. (18) originates from the fact that the fragment  $G$  is exactly spanned by random closed walks of length  $2n$ .

#### 4. Formulas for computing moments

Based on results given above, moments of acyclic chains can be derived one by one by proceeding from the lowest member to higher ones. We list formulas with  $n$  up to 7 below

$$\begin{aligned}
\mu_0^G &= N_{[1]} \\
\mu_2^G &= 2N_{[2]} \\
\mu_4^G &= 2N_{[2]} + 4N_{[3]} \\
\mu_6^G &= 2N_{[2]} + 12N_{[3]} + 6N_{[4]} + 12N_{[31]} \\
\mu_8^G &= 2N_{[2]} + 28N_{[3]} + 32N_{[4]} + 72N_{[31]} + 8N_{[5]} + 16N_{[41]} \\
\mu_{10}^G &= 2N_{[2]} + 60N_{[3]} + 120N_{[4]} + 300N_{[31]} + 60N_{[5]} + 140N_{[41]} \\
&\quad + 10N_{[6]} + 20N_{[51]} + 20N_{[501]} + 40N_{[411]} \\
\mu_{12}^G &= 2N_{[2]} + 124N_{[3]} + 390N_{[4]} + 1080N_{[31]} + 300N_{[5]} + 804N_{[41]} \\
&\quad + 96N_{[6]} + 216N_{[51]} + 228N_{[501]} + 504N_{[411]} \\
&\quad + 12N_{[7]} + 24N_{[61]} + 24N_{[601]} + 48N_{[511]} + 48N_{[5101]} + 24N_{[502]} \\
\mu_{14}^G &= 2N_{[2]} + 252N_{[3]} + 1176N_{[4]} + 3612N_{[31]} + 1260N_{[5]} + 3836N_{[41]} \\
&\quad + 588N_{[6]} + 1484N_{[51]} + 1652N_{[501]} + 3976N_{[411]} \\
&\quad + 140N_{[7]} + 308N_{[61]} + 322N_{[601]} + 728N_{[511]} + 672N_{[5101]} \\
&\quad + 366N_{[502]} \\
&\quad + 14N_{[8]} + 28N_{[71]} + 28N_{[701]} + 28N_{[7001]} + 28N_{[602]} \\
&\quad + 56N_{[611]} + 56N_{[6101]} + 56N_{[61001]} + 56N_{[6011]} + 56N_{[512]} \\
&\quad + 112N_{[5111]}.
\end{aligned}
\tag{19}$$

The first two equalities are obvious, following the definitions of zero- and second-moments. The derivation of other members is also straightforward. We illustrate  $\mu_4^G$  and  $\mu_6^G$  for instance, for which coefficients in Eq. (11) are readily obtained from Eqs. (17) and (18). Eq. (17) gives

$$C_4^{[2]} = 2, \quad C_6^{[2]} = 2, \quad C_6^{[3]} = 12$$

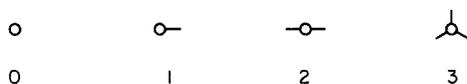
and Eq. (18) gives

$$C_4^{[3]} = 2 \times 2 = 4, \quad C_6^{[4]} = 2 \times 3 = 6, \quad C_6^{[31]} = 2 \times 3 \times 2 = 12.$$

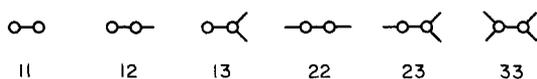
Eq. (19) shows how moments are dependent on connectivities in detail, where not only bonds but also multi-atom fragments play roles. Notice also that smaller fragments appear more frequently than larger ones, which seems to imply the nearer neighbor interactions are more important.

### 5. Alternative description

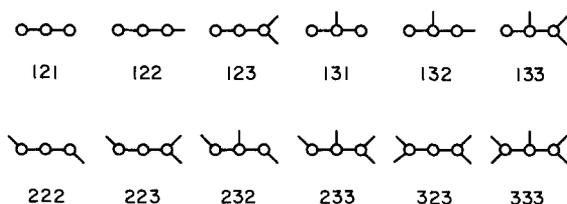
The fragment set used to characterize connectivity is not unique. Various choices are possible. An alternative scheme [15] which is useful in practice is based on the characterization of each atom by its valency (or degree) in detail. Thus, there are four species of atoms, specified by zero-, mono-, bi- and tri-valencies shown diagrammatically below



Similarly, we have six types of bonds characterized by a pair of atomic valencies



and twelve tri-atomic fragments with a set of triple atomic valencies



and so on.

If we deal with one component systems other than free atoms and ethylene, we are able to ignore the trivial fragments 0 and 11, and write down the relations holding for numbers of various fragments in the same set as well as between

different sets. These include

$$\begin{aligned}
 N_1 + N_2 + N_3 &= N_{[1]} \\
 N_1 + 2N_2 + 3N_3 &= 2[N_{12} + N_{13} + N_{22} + N_{23} + N_{33}] = 2N_{[2]} \\
 2N_2 + 6N_3 &= N_{12} + 2N_{13} + 2N_{22} + 3N_{23} + 4N_{33} \\
 &= 2\{N_{121} + N_{123} + N_{131} + N_{132} + N_{133} \\
 &\quad + N_{222} + N_{223} + N_{232} + N_{233} + N_{323} + N_{333}\} = 2N_{[3]} \\
 &\dots
 \end{aligned} \tag{20}$$

and others which we omit for simplicity. By means of Eq. (20) moments in Eq. (19) can be transformed in terms of the new set as follows

$$\begin{aligned}
 \mu_0^G &= \frac{1}{6}\{9N_{12} + 8N_{13} + 6N_{22} + 5N_{23} + 4N_{33}\} \\
 \mu_2^G &= 2N_{12} + 2N_{13} + 2N_{22} + 2N_{23} + 2N_{33} \\
 \mu_4^G &= 4N_{12} + 6N_{13} + 6N_{22} + 8N_{23} + 10N_{33} \\
 \mu_6^G &= 8N_{12} + 18N_{13} + 20N_{22} + 36N_{23} + 58N_{33} \\
 \mu_8^G &= 16N_{12} + 54N_{13} + 62N_{22} + 148N_{23} + 298N_{33} \\
 &\quad + 8N_{222} + 16N_{223} + 32N_{323} + 8N_{232} + 16N_{233} + 32N_{333} \\
 &\dots
 \end{aligned} \tag{21}$$

where we omit moments higher than  $\mu_8^G$ .

These formulas emphasize the role of atomic valency characteristic of each fragment, in addition to that of fragment length. This is well known as the branching effect by chemists. Eq. (21) emphasizes again the dominant role of bonds. They appear uniquely before  $\mu_6^G$  and play a major role in the subsequent moments.

## 6. Total $\pi$ electron energy

Because the pairing theorem holds for alternants, the total  $\pi$  electron energy of the ground state of acyclic systems may be written as

$$E_\pi = 2 \sum_{i=\text{occ}} E_i = \sum_{i=1}^N |E_i| \tag{22}$$

where the latter summation runs through the entire orbital set. The quantity  $|E|$  means the absolute value of  $E$ , which is an even behaved function of  $E$  and can be expanded in terms of even powers of  $E$  [6]

$$|E| = a_0 + a_2 E^2 + a_4 E^4 + \dots \tag{23}$$

provided the point  $E=0$  is excluded (which actually has no effect on the evaluation of  $E_\pi$ ), or empirically approximated by a polynomial of  $E^2$  [7].

On substituting Eq. (23) into Eq. (22) and utilizing the expression for moments, we have

$$E_{\pi} = a_0\mu_0 + a_2\mu_2 + a_4\mu_4 + \dots \quad (24)$$

On further substituting either Eq. (19) or (21) into Eq. (24),  $E_{\pi}$  is transformed to

$$E_{\pi} = \sum_K a_{[K]} N_{[K]} \quad (25)$$

$$E_{\pi} = \sum_{i,j,k,\dots} E_{ijk\dots} N_{[ijk\dots]} \quad (26)$$

where the numbers of various fragments are used as variables. Thus the coefficient  $a_{[K]}$  or  $E_{ijk\dots}$  means the energy component contributed by fragment  $[K]$  or  $ijk\dots$ . These formulas display the desired feature of additivity of the total energy in terms of the constituents of a molecule.

One should be able, in principle, to determine a coefficient set  $\{a_{[K]}\}$  or  $\{E_{ijk\dots}\}$ . A set of known values of  $E_{\pi}$  can be used for fitting Eq. (25) or (26) if they are truncated. For example, the following approximate formula comes from Eq. (26)

$$E_{\pi} = E_{12}N_{12} + E_{13}N_{13} + E_{22}N_{22} + E_{23}N_{23} + E_{33}N_{33} \quad (27)$$

by eliminating higher terms than numbers of bonds. Known  $E_{\pi}$  values of 42 acyclic conjugated molecules have been used for fitting Eq. (27) by a least squares procedure. In this way, the coefficients  $E_{ij}$  are determined, as tabulated in the following table.

**Table 1.** Five  $E_{ij}$  determined by least square procedure

$E_{12}$	$E_{13}$	$E_{22}$	$E_{23}$	$E_{33}$
1.5898	1.4145	1.2691	1.1328	1.0221

A significant trend is shown by this Table in that the bond energy decreases with branching.

## 7. Aromaticity analysis

Based on Eq. (27) and Table 1, a similar analysis of aromaticity of cyclic conjugated systems can be carried out by utilizing the criterion proposed by Dewar [8, 9] as well as by Hess and Schaad [10]. Values of  $E_{\pi}$  of real systems are available from Tables [11] and  $E_{\pi}$  of their corresponding references are calculated from Eq. (27) by enumerating the numbers of the five types of bonds. Then the resonance energy per electron (REPE) is calculated according to

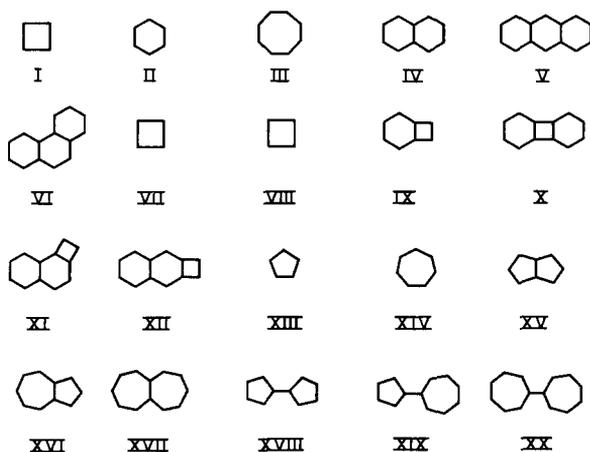
$$\text{REPE} = \frac{1}{N} [E_{\pi} - E_{\pi}(\text{ref.})]. \quad (27)$$

**Table 2.** REPE of some conjugated molecules

Molecule	$E_\pi$	$E_\pi$ (Ref.)	REPE	REPE ( $H-S$ )
I	4.000	5.0764	-0.269	-0.268
II	8.000	7.6146	0.064	0.065
III	9.6569	10.1528	-0.062	
IV	13.6832	13.1679	0.052	0.055
V	19.3137	18.7212	0.042	0.047
VI	19.4531	18.7468	0.050	0.055
VII	7.2077	7.3858	-0.030	-0.028
VIII	9.6568	9.7464	-0.011	-0.010
IX	10.3812	10.6297	-0.031	-0.027
X	16.5055	16.2342	0.023	0.027
XI	15.9962	16.2086	-0.018	-0.012
XII	16.2010	16.1830	0.002	0.007
XIII	7.4659	7.4874	-0.004	-0.002
XIV	9.9943	10.0256	-0.004	-0.002
XV	10.4556	10.6297	-0.022	-0.018
XVI	13.3635	13.1679	0.020	0.023
XVIII	12.7992	13.1679	-0.037	-0.033
XIX	15.9306	15.7061	0.019	0.022
XX	18.0048	18.2443	-0.017	-0.014

REPE's of acyclic systems calculated in this way vary from  $-0.005$  to  $0.005$  ( $\beta$ ), which gives the range of non-aromaticity. Compounds with their REPE outside this range should be aromatic (+) or antiaromatic (-).

Dozens of compounds have been calculated, and the result seems parallel to that of Hess and Schaad [10]. Only part of our results is listed in Table 2, which includes the following compounds



A few points seem worthwhile emphasizing: (1) The number of parameters used here is limited to five possible bonds. They are uniquely determined, without arbitrariness, and the result (including a large number of compounds) may be

refined upon including more parameters. (2) The numbers of various types of bonds are also uniquely enumerable, without averaging in terms of Kekule structures. So, it seems available for non-classical systems, too. We feel the moment analysis gives a sound basis for understanding additivity and aromaticity. It should be extendable to other chemical problems [12].

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#### References

1. See, for example Dewar, M. J. S.: The molecular orbital theory of organic chemistry. New York: McGraw-Hill 1969
2. Graovac, A., Gutman, I, Trinajstic, N.: Topological approach to the chemistry of conjugated molecules. Berlin, Heidelberg, New York: Springer-Verlag 1977
3. Tang, Au-chin, Kiang, Yuan-sun: *Sci. Sin.* **19**, 208 (1976)
4. Kiang Yuan-sun: *Int. J. Quantum Chem.* **28**, 331 (1980); *Int. J. Quantum Chem.* **S15**, 293 (1981)
5. Cao, Wei-liang, Jiang, Yuan-sheng: *Acta Chimica Sinica* **40**, 880 (1982)
6. Gutman, I, Trinajstic, N.: *Chem. Phys. Lett.* **17**, 535 (1972)
7. Gutman, I, Trinajstic, N.: *Chem. Phys. Lett.* **20**, 257 (1973)
8. Dewar, M. J. S., Gleicher, G. J.: *J. Am. Chem. Soc.* **87**, 685, 692 (1965)
9. Dewar, M. J. S., and de Llano, C.: *J. Am. Chem. Soc.* **91**, 789 (1969)
10. Hess, Jr., B. A., Schaad, L. J.: *J. Am. Chem. Soc.* **93**, 305 (1971)
11. See, for example E. Heilbronner, P. A. Straub; "HMO" Hückel molecular orbitals. New York: Springer-Verlag, 1966
12. Burdett, J. K., Lee, S., Sha, W. C.: to be published in *Croat. Chem. Acta*.

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